

# Two-dimensional elastic Bloch waves in helical periodic structures

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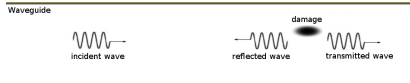
# Context

## Guided wave applications:

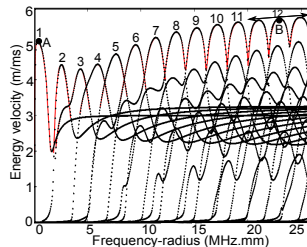
- dynamic analysis of elongated structures
- propagation over long distances, sensitivity to small damages
- examples:
  - Non Destructive Evaluation (ultrasonics)
  - vibration and noise reduction
  - statistical energy analysis...

## Generality about waveguides:

- Guided wave propagation: dispersive and multimodal
- Dispersion curves required
- Modeling tools needed



**NDE:** detecting a damage with elastic waves



Energy velocity vs. frequency in a cylindrical bar

# On the modeling of elastic waveguides

## Full 3D approach:

- high frequency (e.g. ultrasonics) → fine mesh
- guided waves go to infinity → large model
- huge computational memory required
- tedious post-processing for wave modes...

## Reduced modeling (modal approach):

- guided waves = modes
- eigenvalue problem
- plates, cylinders: analytical approaches (Thomson-Haskell, ©Disperse, ...)
- **arbitrary geometry**: finite element discretization
  - of cross-section<sup>1</sup> (often referred to as “SAFE” method)
  - of a 3D slice with Bloch-Floquet periodic conditions<sup>2</sup> (“WFEM”)

SAFE: Semi-Analytical Finite Element method

**WFEM: Wave Finite Element Method**

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<sup>1</sup>Lagasse JASA 1973, Aalami JAM 1973, Hayashi et al. Ultrasonics 2003, Bartoli et al. JSV 2006,...

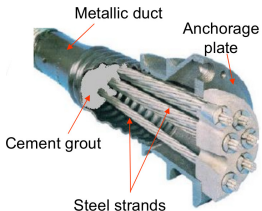
<sup>2</sup>Gry et al. JSV 1997, Mace et al. JASA 2005,...

# Our motivation: helical structures

**Typical example of application:** NDT and SHM of cables (damage detection, tension estimation,...)

## Helical symmetry:

- along one direction (beam-like structures) or two directions (tube-like)
- continuous (uniform waveguides) or discrete (periodic waveguides)

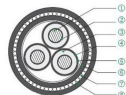


**Left:** bridge cable (anchorage), **right:** seven-wire strand



Construction

- ① Conductor: Compact stranded copper conductor, Cl. 2 as per IEC 60228
- ② Conductor Screen: Semi-conductor
- ③ Insulation: XLPE (cross-linked polyethylene) rated at 90°C
- ④ Insulation Screen: Semi-conductor
- ⑤ Screen: Copper tape
- ⑥ Inner arminging - PVC
- ⑦ Arminging: Galvanized steel wire - 3 cores
- ⑧ Sheath: PVC or FR-PVC type ST2 to IEC 60602, black

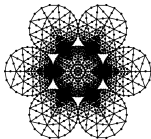


**Umbilical power cable**

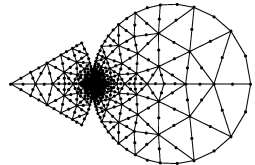
# State-of-the-art (GeoEND lab)



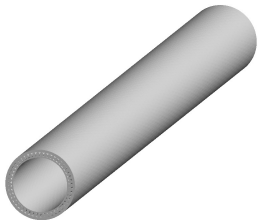
From a 3D seven-wire strand...



... to SAFE 2D  
(continuous helical symmetry,  
12369 dofs)



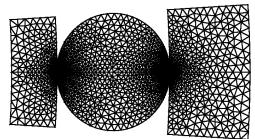
... and to SAFE 2D/6  
(+discrete rot. symmetry,  
2094 dof → CPU time/13)



From a 3D cable armor...



... to SAFE 2D  
(1,000,000 dofs)

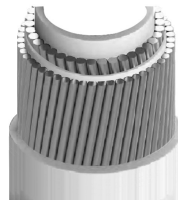


... and to 2D/50  
(22587 dofs)

# Let us further generalize...



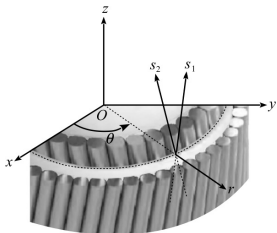
Doubled armored cable  
(high strength submarine  
power cable)  
source: BPP catalog



Double armor

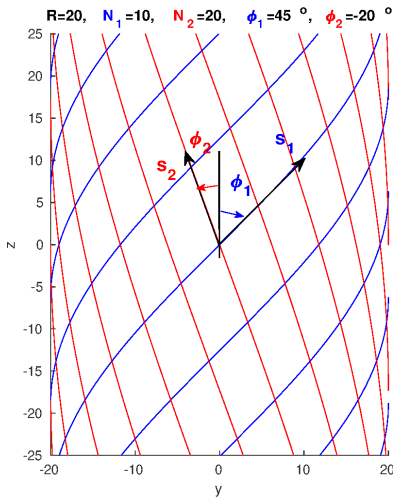
**Discrete** symmetry along **two directions** ? (both helical)

# Helical coordinates



Cartesian, cylindrical and helical coordinates

- Two layers (virtual or not), Layer 1 (blue) and Layer 2 (red)
- Layer  $\alpha$  is divided by  $N_\alpha$  helices, oriented along  $s_\alpha$ , of radius  $R_\alpha$  and lay angle  $\phi_\alpha$  ( $\alpha=1,2$ )



Bi-helical periodicity pattern  
 $(s_1, s_2)$ : helical coordinates

## Bi-helical coordinate system

- Relationship between cylindrical and helical coordinates:

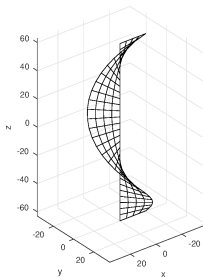
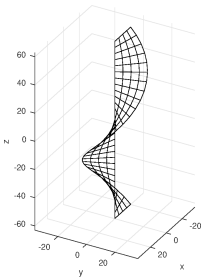
$$\begin{cases} \theta = \frac{2\pi}{\ell_1} s_1 + \frac{2\pi}{\ell_2} s_2 & L_{1,2} : \text{helix steps measured along z-axis (straight)} \\ z = \frac{L_1}{\ell_1} s_1 + \frac{L_2}{\ell_2} s_2 & \ell_{1,2} : \text{curvilinear steps} \end{cases}$$

- Position vector  $\mathbf{OM} = r\mathbf{e}_r(\theta) + z\mathbf{e}_z$ :

$$\mathbf{OM} = r \cos\left(\frac{2\pi}{\ell_1} s_1 + \frac{2\pi}{\ell_2} s_2\right) \mathbf{e}_x + r \sin\left(\frac{2\pi}{\ell_1} s_1 + \frac{2\pi}{\ell_2} s_2\right) \mathbf{e}_y + \left(\frac{L_1}{\ell_1} s_1 + \frac{L_2}{\ell_2} s_2\right) \mathbf{e}_z$$

The 3D coordinates are now:

$$(s_1, s_2, r)$$



Surfaces  $s_2=\text{cst}$  (left),  $s_1=\text{cst}$  (right)



# Unit cell boundaries

## A first difficulty...

Left/right boundaries  $\Gamma_1^\pm$ :

$$s_1 = \text{cst} = \pm \Delta \ell_1 / 2 \text{ (helicoids)}$$

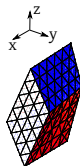
Bottom/top boundaries  $\Gamma_2^\pm$ :

$$s_2 = \text{cst} = \pm \Delta \ell_2 / 2 \text{ (helicoids)}$$

with cell lengths defined by:

$$\Delta \ell_1 = \frac{\ell_1}{N_2} \frac{L_2}{L_2 - L_1}, \Delta \ell_2 = \frac{\ell_2}{N_1} \frac{L_1}{L_2 - L_1}$$

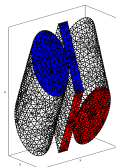
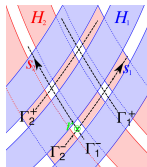
→ **Unit cells are cut by non-plane surfaces**



Unit cell for a thick hollow cylinder

## Parametrization of boundaries for a double armor:

- intersection between 4 helicoids  $\Gamma_\alpha^\pm$  (2D) and the helical wire volumes  $H_\alpha$  (3D)...
- the solution is not analytic  
→ multiple non-linear equations to solve (in one variable)



Unit cell for a double armor  
left: front view, right: FE mesh

## Invariance(\*) in 2D periodic media

Reminder: Bloch waves in two directions  $s_1$  and  $s_2$ :  $\psi(s_1, s_2) = e^{ik_1 s_1} e^{ik_2 s_2} u(s_1, s_2)$ , with  $k_\alpha$ : wavenumbers,  $u$ :  $\Delta\ell_\alpha$ -periodic function ( $\Delta\ell_\alpha$ : unit cell lengths)

**Question: existence of Bloch waves in such a geometry ?**

- The coeff. of diff. equations must be periodic in two curved directions
- Let us rewrite the equilibrium equations in the bi-helical system...

Example: equilibrium equation for elasticity

$$\sigma_{,j}^{ij} + \Gamma_{mj}^i \sigma^{mj} + \Gamma_{mj}^j \sigma^{im} + \rho\omega^2 g^{ij} u_j = 0, \text{ with } \sigma^{ij} = C^{ijkl} \epsilon_{kl}, \epsilon_{kl} = \frac{1}{2}(u_{k,l} + u_{l,k} - \Gamma_{kl}^m u_m)$$

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**The coefficients are functions of:** the physical properties ( $\rho$  and  $C^{ijkl}$ )... and  $\Gamma_{ij}^k$ !

**Key formula:**  $\Gamma_{ij}^k = \frac{1}{2} g^{kl} \left( \frac{\partial g_{jl}}{\partial x^i} + \frac{\partial g_{il}}{\partial x^j} - \frac{\partial g_{ij}}{\partial x^l} \right) \rightarrow$  functions of the  $g_{ij}$ 's only

**Conditions for invariance in periodic media along curved directions ( $s_1, s_2$ ):**

- 1 the shape of the geometry is  $(\Delta\ell_1, \Delta\ell_2)$ -periodic along  $(s_1, s_2)$
- 2 the physical properties are  $(\Delta\ell_1, \Delta\ell_2)$ -periodic along  $(s_1, s_2)$
- 3 **the metric tensor  $g$  does not depend on  $(s_1, s_2)$**

(\*) "invariance"  $\equiv$  the coefficients of the differential equations are periodic

# Proof of existence in bi-helical structures

Application to the bi-helical coordinates: let us calculate the metric tensor...

Reminder:  $(\mathbf{g})_{ij} = \mathbf{g}_i \cdot \mathbf{g}_j$ , with 
$$\begin{cases} \mathbf{g}_1 = \frac{\partial \mathbf{OM}}{\partial s_1} \\ \mathbf{g}_2 = \frac{\partial \mathbf{OM}}{\partial s_2} \\ \mathbf{g}_3 = \frac{\partial \mathbf{OM}}{\partial r} \end{cases}$$

$$\Rightarrow \mathbf{g} = \begin{bmatrix} \frac{4\pi^2 r^2 + L_1^2}{\ell_1^2} & \frac{4\pi^2 r^2 + L_1 L_2}{\ell_1 \ell_2} & 0 \\ \frac{4\pi^2 r^2 + L_1 L_2}{\ell_1 \ell_2} & \frac{4\pi^2 r^2 + L_2^2}{\ell_2^2} & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad \text{point 3 checked } \checkmark$$

Corollary: proof of the existence of Bloch waves in bi-helical waveguides whatever the physics

# Implementation: the Wave Finite Element Method

FE discretization of the unit cell:

$$(\mathbf{K} - \omega^2 \mathbf{M} - i\omega \mathbf{C})\mathbf{U} = \mathbf{F}$$

Elastic fields  $\mathbf{U}$  and  $\mathbf{F}$  are not scalar: Cartesian components must be transformed **in the covariant basis**

$$\begin{bmatrix} \mathbf{g}_1 \\ \mathbf{g}_2 \\ \mathbf{g}_3 \end{bmatrix} = \mathbb{J}^T \begin{bmatrix} \mathbf{e}_x \\ \mathbf{e}_y \\ \mathbf{e}_z \end{bmatrix}, \quad \mathbb{J}^T = \begin{bmatrix} -\frac{2\pi r}{\ell_1} \sin \theta & \frac{2\pi r}{\ell_1} \cos \theta & \frac{L_1}{\ell_1} \\ -\frac{2\pi r}{\ell_2} \sin \theta & \frac{2\pi r}{\ell_2} \cos \theta & \frac{L_2}{\ell_2} \\ \cos \theta & \sin \theta & 0 \end{bmatrix}$$

Apply:

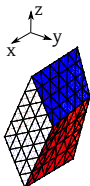
- 1 Displacement boundary conditions of Floquet-Bloch type (two-directional):

$$\mathbf{J}_R^T \mathbf{U}_R = \lambda_1 \mathbf{J}_L^T \mathbf{U}_L, \quad \mathbf{J}_T^T \mathbf{U}_T = \lambda_2 \mathbf{J}_B^T \mathbf{U}_B, \quad \mathbf{J}_{RB}^T \mathbf{U}_{RB} = \lambda_1 \mathbf{J}_{LB}^T \mathbf{U}_{LB}$$

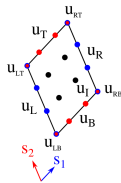
$$\mathbf{J}_{LT}^T \mathbf{U}_{LT} = \lambda_2 \mathbf{J}_{LB}^T \mathbf{U}_{LB}, \quad \mathbf{J}_{RT}^T \mathbf{U}_{RT} = \lambda_1 \lambda_2 \mathbf{J}_{LB}^T \mathbf{U}_{LB}$$

with:  $\lambda_1 = e^{ik_1 \Delta \ell_1}$ ,  $\lambda_2 = e^{ik_2 \Delta \ell_2}$  ( $k_1, k_2$ : helical wavenumbers)

- 2 Force boundary conditions... (by condensation of displacement)
- 3 Solve the eigenvalue problem for  $\omega$  ( $k_{1,2}$  are the fixed parameters)



Unit cell example  
(a thick tube)



dofs  
classification  
(2D view)

## Relationship between wavenumbers

- Our structure is of cylindrical type:

$$\begin{cases} k_z = \frac{k_2 \ell_2 - k_1 \ell_1}{L_2 - L_1} \\ k_\theta = \frac{k_1 \ell_1 L_2 - k_2 \ell_2 L_1}{2\pi(L_2 - L_1)} \end{cases}$$

- As opposed to a plate-like geometry, the propagation constants  $\lambda_1$  and  $\lambda_2$  are not independent:

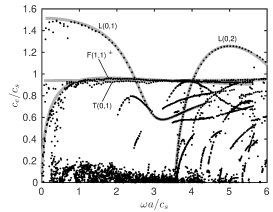
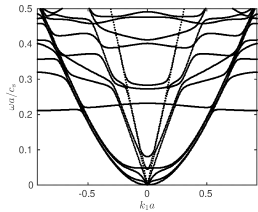
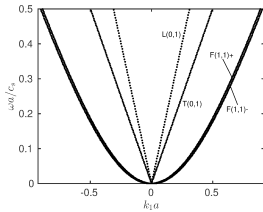
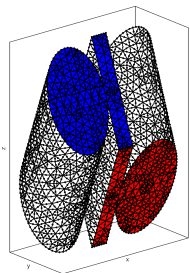
$$\lambda_1^{N_2} \lambda_2^{-N_1} = 1$$

⇒ For  $(k_1, k_2)$  in the first Brillouin zone, the circumferential wavenumber is given by:

$$k_\theta = n_\alpha$$

with  $n_\alpha$  varying between  $N_1$  consecutive integer values or  $N_2$  consecutive values depending on whether the user sets  $k_1$  or  $k_2$

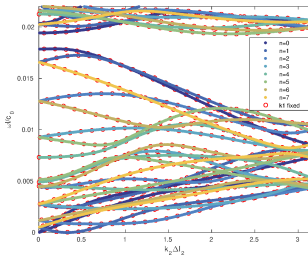
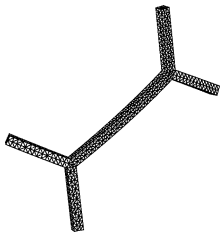
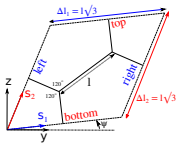
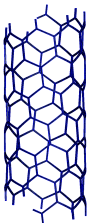
# Application to double armor cables



Dispersion curves for: single free wire (left), double armor in low-frequency regime (middle), double armor in high-frequency regime (right, gray: free wire) computed for  $n_1 = 0$ .

## Chiral nanotubes

Test case taken from from Maurin et al., CMAME, 2017



Left: FE mesh of the unit cell. Right: dispersion curves



# Conclusion and future works

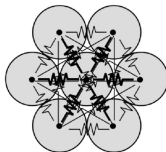
## Conclusion and short-term works:

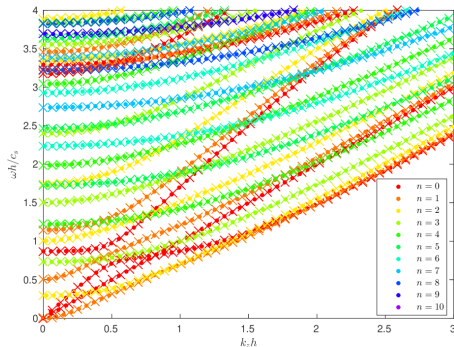
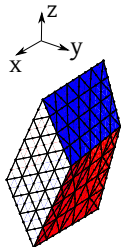
- a modeling framework for Bloch waves in bi-helical structures
- a straightforward extension: beam-like homogenization (quasi-static cable bending fatigue)
- another bi-helical architecture: the double layer strand



## What's next in the mid-term?

- considering more and more heterogeneity...
- work in progress (with Pierric Mora)...
- ... cables simplified as “granular” media (contact theory + beam theory)





**Thick tube test case ( $R/h = 2$ ).** Left: FE mesh of the bi-helical unit cell, right: dispersion curves for various circumferential orders  $n$ , computed with the bi-helical unit cell ( $\bullet$ ), with a commonly used straight unit cell ( $\times$ ) cut along  $\theta$  and  $z$  ( $N_\theta = 40$ ,  $\Delta l_z = 0.5h$ )

## Invariance in arbitrarily curved structures ?

**Invariance along  $s$ :** the coefficients of equilibrium equations, including boundary conditions, must be  $s$ -periodic...

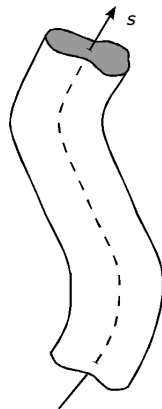
Do guided waves exist in curved structures? Not always, obviously...

Conditions for invariance in a curved direction  $s$ :

- ① cross-section ✓
- ② physical properties ✓
- ③ ?

A third condition is required:

- intuition: difficult...
- inherent to the coordinate system considered (curvilinear)
- answer: **differential geometry and tensor analysis** required...



# Elements of differential geometry

## Definitions:

- covariant basis:  $(\mathbf{g}_1, \mathbf{g}_2, \mathbf{g}_3) = (\partial\mathbf{X}/\partial x, \partial\mathbf{X}/\partial y, \partial\mathbf{X}/\partial s)$
- contravariant basis  $\mathbf{g}^i$ :  $\mathbf{g}_i \cdot \mathbf{g}^j = \delta_i^j$
- Christoffel symbols  $\Gamma_{ij}^k$ :  $\mathbf{g}_{i,j} = \Gamma_{ij}^k \mathbf{g}_k$  ( $\Leftrightarrow \Gamma_{ij}^k = \mathbf{g}_{i,j} \cdot \mathbf{g}^k$ )

The coefficients of differential operators are given by the  $\Gamma_{ij}^k$ 's

## Examples:

$$\nabla\varphi = \varphi_{,i} \mathbf{g}^i$$

$$\nabla\mathbf{u} = (u_{i,j} - \Gamma_{ij}^k u_k) \mathbf{g}^i \otimes \mathbf{g}^j$$

$$\nabla \cdot \boldsymbol{\sigma} = (\sigma_{ij,j} - \Gamma_{ij}^k \sigma_{kj} - \Gamma_{ij}^k \sigma_{ik}) \mathbf{g}^i$$

...

## The metric tensor

- Definition:  $(\mathbf{g})_{ij} = \mathbf{g}_i \cdot \mathbf{g}_j$
- Key formula:  $\Gamma_{ij}^k = \frac{1}{2} g^{kl} \left( \frac{\partial g_{jl}}{\partial x^i} + \frac{\partial g_{il}}{\partial x^j} - \frac{\partial g_{ij}}{\partial x^l} \right) \rightarrow$  functions of the  $g_{ij}$ 's only

### Metric tensor for usual coordinate systems

cartesian:

$$\mathbf{g} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

cylindrical:

$$\mathbf{g} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & r^2 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

spherical:

$$\mathbf{g} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & r^2 & 0 \\ 0 & 0 & r^2 \sin^2 \varphi \end{bmatrix}$$

### Metric tensor for some helical systems

- one-directional helical  $(x, y, s)$ :

$$\mathbf{X}(x, y, s) = \mathbf{R}(s) + x\mathbf{N}(s) + y\mathbf{B}(s)$$

$$\mathbf{g} = \begin{bmatrix} 1 & 0 & -\tau y \\ 0 & 1 & \tau x \\ -\tau y & \tau x & \tau^2(x^2 + y^2) + (1 - \kappa x)^2 \end{bmatrix}$$

- mixed helical-polar  $(\rho, \theta, s)$ :

$$\mathbf{X}(\rho, \theta, s) = \mathbf{R}(s) + \rho \cos \theta \mathbf{N}(s) + \rho \sin \theta \mathbf{B}(s)$$

$$\mathbf{g} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & r^2 & \tau^2 r^2 \\ 0 & \tau^2 r^2 & \tau^2 r^2 + (1 + \kappa r \cos \theta)^2 \end{bmatrix}$$

